

# The Influence of Inclined Magnetic Field and Heat Transfer on the MHD Convective Flow in a Vertical Channel Filled Partly with Porous Medium

P. T. Hemamalini, M. Shanthi

**ABSTRACT:** An analytic study has been made of a laminar fully developed MHD flow bounded by infinite vertical parallel plates with effect of inclined magnetic field partly filled with fluid and partly with porous matrix. The motions of the plates are in the opposite direction and are maintained at distinct temperatures. The perturbation method has been chosen to derive the expression for velocity flow and temperature distribution and the effect on flow velocity and temperature due to magnetic field and Darcy number has been illustrated for the fluid and porous region with the help of graph.

**KEYWORDS:** MHD Convection Flow, Inclined Magnetic Field, Vertical Channel, Darcy number, partly filled porous medium.

## I. INTRODUCTION

A significant interest has been created by recent technological implications in the study of flow in channels through partly filled porous medium. An extensive research has been carried out in the multiphase flow through channels and different wall temperatures in the presence of inclined magnetic field. Due to its common occurrence in engineering fields such as in thermal energy storage system, petrol purification, etc., it would be interesting to study the MHD flow in a region with partly loaded with porous medium in the presence of inclined magnetic field.

A. K. Singh, and T. Paul (2006) considered transient natural convection between two vertical walls heated/cooled asymmetrically. M. M. S. El-Din (2007) investigated on the fully developed mixed convection laminar flow with viscous dissipation in a uniformly heated vertical double-passage channel. Radiation effects on an unsteady MHD convective heat and mass transfer flow past a semi-infinite vertical permeable moving plate embedded in a porous medium was illustrated by V. R. Prasad and N. B. Reddy, (2008). D. S. Chauhan and V. Kumar, (2009) investigated on the effects of slip conditions on forced convection and entropy generation in a circular channel occupied by a highly porous medium: Darcy extended Brinkman-Forchheimer model. Srivastava and Singh. A. K., (2010) considered the mixed convection in a composite system bounded by vertical walls.

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Free convection flow through vertical plates moving in different direction and partially filled with porous medium was discussed by Umesh Gupta, et al, (2011). Magnetohydrodynamic convection effects with viscous and ohmic dissipation in a vertical channel partially filled by a porous medium was illustrated by Dileep Singh Chauhan and Rashmi Agrawal, (2012). Prasad, K. V., K. Vajravelu, et al, (2013) studied MHD flow with heat transfer in a power-law liquid film at a porous surface in the presence of thermal radiation. S. Purkayastha and Choudhary .R (2014) considered visco-elastic effects on convection flow through a vertical rotating channel partially filled with a porous medium. Adeyemi Isaiyah Fagbade, et al (2015) observed the influence of magnetic field, viscous dissipation and thermophoresis on Darcy-Forchheimer mixed convection flow in fluid saturated porous media. Free convection flow of a Jeffrey fluid between vertical plates partially filled with porous medium was learnt by S. Dhananjaya1, et al (2015), A. Adeniyani, et al, (2016) discussed mixed convection radiating flow and heat transfer in a vertical channel partially filled with a Darcy-Forchheimer Porous Substrate. V. G. Gupta, Ajay Jain, et al., (2016) studied the convective effects on MHD flow and heat transfer between vertical plates moving in opposite direction and partially filled with a porous medium. Heat, Mass Transfer in free Convective Flow of Walter's Liquid Model-B through Rotating Vertical Channel was illustrated by Pooja Sharma, et al. (2017).

The aim of the study is to analyze the fully developed laminar flow in the vertical channel partially porous medium and clear fluid in presence of inclined magnetic field.

## II. PHYSICAL DESCRIPTION OF THE PROBLEM

A vertical channel of width  $H$ , partly loaded with clear fluid and partially with porous matrix having interface, is considered with inclined magnetic field, when one of the plates is heated and the other is cooled. The fully developed laminar free-convection flow confined in the vertical channel is illustrated. The Cartesian co-ordinate system is so taken that  $X$  axis is marked along one of the plates but opposite to the gravitational field and  $Y$  axis perpendicular to it. The vertical plates are placed at  $[\bar{Y}=0]$ , and  $[\bar{Y}=H]$  are moving in opposite directions with the velocities  $\bar{U}_f$  and  $\bar{U}_p$  in the  $X$  direction. Let the wall temperature at  $\bar{Y}=0$  be  $\bar{T}_f = [\bar{T}_c + A(\bar{T}_h - \bar{T}_c)]$  and

$\bar{Y} = H$  be  $\bar{T}_p = [\bar{T}_c + B(\bar{T}_h - \bar{T}_c)]$  respectively.

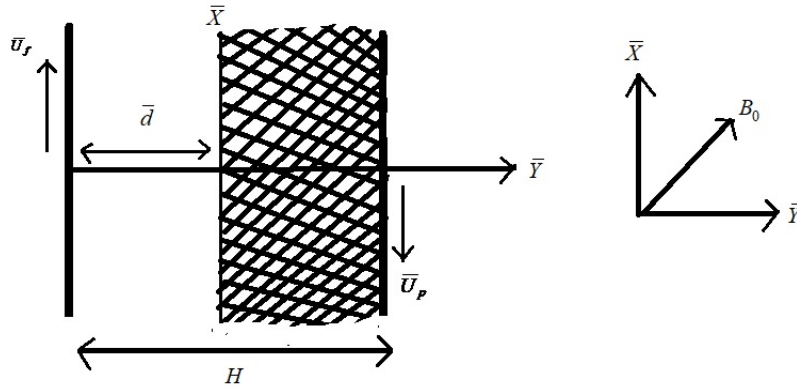


Figure 1 Geometric model and coordinate system

### III. BASIC EQUATIONS AND METHODOLOGY

The governing equations for the above physical configuration are illustrated by:  
Clear Fluid Region

$$\frac{\partial^2 \bar{U}_f}{\partial \bar{y}^2} + \frac{g\beta}{\nu} (\bar{T}_f - \bar{T}_c) - \frac{B_0 \sigma}{\rho \nu} \sin \theta \bar{U}_f = 0$$

(1)

$$\left( \frac{d\bar{U}_f}{d\bar{y}} \right)^2 + \frac{k}{\nu} \frac{d^2}{d\bar{y}^2} (\bar{T}_f - \bar{T}_c) = 0 \quad (2)$$

Porous fluid Region

$$\frac{\partial^2 \bar{U}_p}{\partial \bar{y}^2} + \frac{g\beta}{\nu} (\bar{T}_p - \bar{T}_c) - \frac{B_0 \sigma}{\rho \nu} \sin \theta \bar{U}_p - \frac{\bar{U}_p}{\bar{k}} = 0 \quad (3)$$

$$\left( \frac{d\bar{U}_p}{d\bar{y}} \right)^2 + \frac{k}{\nu} \frac{d^2}{d\bar{y}^2} (\bar{T}_p - \bar{T}_c) - \frac{\bar{U}_p^2}{\bar{k}} = 0 \quad (4)$$

The initial, boundary conditions are:

$$\begin{aligned} \bar{U}_f &= \frac{g\beta H^2 [\bar{T}_h - \bar{T}_c]}{\nu} u_0; \bar{T}_f = [\bar{T}_c + A(\bar{T}_h - \bar{T}_c)] & \text{at } \bar{y} = 0 \\ \bar{U}_p &= -\frac{g\beta [\bar{T}_h - \bar{T}_c] H^2}{\nu} u_0; \bar{T}_p = [\bar{T}_c + B(\bar{T}_h - \bar{T}_c)] & \text{at } \bar{y} = H \\ \frac{d\bar{U}_f}{d\bar{y}} &= \frac{d\bar{U}_p}{d\bar{y}}; \bar{U}_f = \bar{U}_p; \frac{d\bar{T}_f}{d\bar{y}} = \frac{d\bar{T}_p}{d\bar{y}} = \bar{T}_p; & \text{at } \bar{y} = \bar{d} \end{aligned} \quad (5)$$

Where  $\bar{U}_f$  -denotes the velocity of the plate at ( $Y=0$ ),  $\bar{U}_p$  plate velocity at ( $Y=H$ ),  $B_0$  strength of uniform magnetic field,  $\rho$  density and  $k$  thermal conductivity. The thermal expansion co-efficient of the fluid is marked by  $\beta$ , thermal Grashof number by  $G_r$ , temperature of the heated plate by  $\bar{T}_h$ , temperature of the cooled plate by  $\bar{T}_c$  and the acceleration due to gravity by  $g$ ,  $A$  and  $B$  are the interchangeable temperatures at the plates.  $\sigma$  denotes the electrical conductivity, the permeability of porous medium is marked by  $\bar{k}$  and  $K$  is the thermal conductivity and  $N$  is the buoyancy parameter. Kinematic viscosity and the dynamic viscosity are marked by  $\mu$  and  $\nu$  respectively.

Introducing non-dimensional quantities,

$$\begin{aligned} U_f &= \frac{\nu \bar{U}_f}{g\beta H^2 (\bar{T}_h - \bar{T}_c)}, \theta_f = \frac{(\bar{T}_f - \bar{T}_c)}{(\bar{T}_h - \bar{T}_c)}; Da = \frac{\bar{k}}{H^2}; y = \frac{\bar{y}}{H}; d = \frac{\bar{d}}{H}; \\ U_p &= \frac{\nu \bar{U}_p}{g\beta H^2 (\bar{T}_h - \bar{T}_c)}, \theta_p = \frac{(\bar{T}_p - \bar{T}_c)}{(\bar{T}_h - \bar{T}_c)}; N = \frac{g^2 \beta^2 H^4 (\bar{T}_h - \bar{T}_c)}{k\nu}; M = \frac{\sigma B_0^2 h^2}{\nu} \end{aligned} \quad (6)$$

With reference to the problem under consideration, the basic equations in non-dimensional quantities for the equations (1) to (5) can be expressed as:

Fluid region



$$\frac{d^2U_f}{dy^2} + \theta_f - MU_f = 0 \quad (7)$$

$$\frac{d^2\theta_f}{dy^2} + N\left(\frac{dU_f}{dy}\right)^2 = 0 \quad (8)$$

Porous Region

$$\frac{d^2U_p}{dy^2} + \theta_p - \left(M + \frac{1}{Da}\right)U_p = 0 \quad (9)$$

$$\frac{d^2\theta_p}{dy^2} + N\left(\frac{dU_p}{dy}\right)^2 + \frac{N}{Da}U_p^2 = 0 \quad (10)$$

The acquired dimensionless forms of the boundary conditions are:

$$\begin{aligned} U_f = u_0; \theta_f = A & \quad \text{at} \quad y = 0 \\ U_p = -u_0; \theta_f = B & \quad \text{at} \quad y = 1 \\ \frac{dU_f}{dy} = \frac{dU_p}{dy}; U_f = U_p & \quad \text{at} \quad y = d \\ \frac{d\theta_f}{dy} = \frac{d\theta_p}{dy}; \theta_f = \theta_p & \quad \text{at} \quad y = d \end{aligned} \quad (11)$$

#### IV. SOLUTION METHODOLOGY

The above problem considered is non-linear and can be solved by perturbation technique. Most often in practical problems,  $N$  is assumed to be small. Consider the expansion:

$$\begin{aligned} U_f = U_{0f} + NU_{1f} + O(N^2) & \quad \theta_f = \theta_{0f} + N\theta_{1f} + O(N^2) \\ U_p = U_{0p} + NU_{1p} + O(N^2) & \quad \theta_p = \theta_{0p} + N\theta_{1p} + O(N^2) \end{aligned} \quad (12)$$

Introducing (12) in equation (7) to equation (10), we get

$$\frac{d^2U_{0f}}{dy^2} - MU_{0f} + \theta_{0f} = 0 \quad (13)$$

$$\frac{d^2U_{1f}}{dy^2} - MU_{1f} + \theta_{1f} = 0 \quad (14)$$

$$\frac{d^2\theta_{0f}}{dy^2} = 0 \quad (15)$$

$$\frac{d^2\theta_{1f}}{dy^2} + \left(\frac{dU_{0f}}{dy}\right)^2 = 0 \quad (16)$$

$$\frac{d^2U_{0p}}{dy^2} - \left(\frac{1}{Da} + M\right)U_{0p} + \theta_{0p} = 0 \quad (17)$$

$$\frac{d^2U_{1p}}{dy^2} - \left(\frac{1}{Da} + M\right)U_{1p} + \theta_{1p} = 0 \quad (18)$$

$$\frac{d^2\theta_{0p}}{dy^2} = 0 \quad (19)$$

$$\frac{d^2\theta_{1p}}{dy^2} + \left(\frac{dU_{0p}}{dy}\right)^2 + \frac{1}{Da}U_{0p}^2 = 0 \quad (20)$$

The final boundary conditions are:

$$U_{1f}=0; U_{0f}=u_0; \theta_{1f}=0; \theta_{0f}=A \quad \text{at} \quad y=0:$$

$$U_{1p}=0; U_{0p}=-u_0; \theta_{1p}=0; \theta_{0p}=B; \quad \text{at} \quad y=1:$$

$$U_{1f}=U_{1p}; U_{0f}=U_{0p}; \frac{dU_{1f}}{dy}=\frac{dU_{1p}}{dy}; \frac{dU_{0f}}{dy}=\frac{dU_{0p}}{dy} \quad \text{at} \quad y=d; \quad (21)$$

$$\theta_{1f}=\theta_{1p}; \theta_{0f}=\theta_{0p}; \frac{d\theta_{0f}}{dy}=\frac{d\theta_{0p}}{dy}; \frac{d\theta_{1f}}{dy}=\frac{d\theta_{1p}}{dy} \quad \text{at} \quad y=d$$

The solution for equations (13) - (20) are obtained by using the boundary conditions (21).

$$U_{0f}=a_1e^{\sqrt{M}y}+a_2e^{-\sqrt{M}y}+\frac{A+(B-A)y}{M} \quad (22)$$

$$U_{0p}=a_5e^{\sqrt{K}y}+a_6e^{-\sqrt{K}y}+\left(\frac{A+(B-A)y}{K}\right) \quad (23)$$

$$\theta_{0f}=\theta_{0f}=(A)+(B-A)y \quad (24)$$

$$\theta_{1f}=\frac{-a_1^2e^{2\sqrt{M}y}}{4}+\frac{a_2^2e^{-2\sqrt{M}y}}{4}-\frac{(B-A)^2y^2}{2M^2}+a_1a_2My^2+\frac{2a_2e^{-\sqrt{M}y}}{M\sqrt{M}}(B-A)-\frac{2a_1e^{\sqrt{M}y}}{M\sqrt{M}}(B-A)-K_1y-K_2 \quad (25)$$

$$\theta_{1p}=\frac{a_5^2K_3e^{2\sqrt{K}y}}{4K}+\frac{a_6^2K_3e^{-2\sqrt{K}y}}{4K}+\frac{K_4y}{K^2}+2a_5a_6K_3\frac{y^2}{2}+\frac{A+(B-A)y}{DaK}\frac{2(a_5+a_6)e^{-\sqrt{K}y}}{K}+\frac{(B-A)}{K}\frac{2(a_5+a_6)e^{\sqrt{K}y}}{\sqrt{K}}+K_5y+K_6 \quad (26)$$

$$U_{1f}=a_3e^{\sqrt{M}y}+a_4e^{-\sqrt{M}y}+\frac{1}{M}\left[\frac{-a_1^2e^{2\sqrt{M}y}}{4}+\frac{a_2^2e^{-2\sqrt{M}y}}{4}-\frac{(B-A)^2y^2}{2M^2}+a_1a_2My^2+\frac{2a_2e^{-\sqrt{M}y}}{M\sqrt{M}}(B-A)-\frac{2a_1e^{\sqrt{M}y}}{M\sqrt{M}}(B-A)-K_1y-K_2\right] \quad (27)$$

$$U_{1p}=a_7e^{\sqrt{K}y}+a_8e^{-\sqrt{K}y}+\frac{1}{K}\left[\frac{a_5^2K_3e^{2\sqrt{K}y}}{4K}+\frac{a_6^2K_3e^{-2\sqrt{K}y}}{4K}+\frac{K_4y}{K^2}+2a_5a_6K_3\frac{y^2}{2}+\frac{A+(B-A)y}{DaK}\frac{2(a_5+a_6)e^{-\sqrt{K}y}}{K}+\frac{(B-A)}{K}\frac{2(a_5+a_6)e^{\sqrt{K}y}}{\sqrt{K}}+K_5y+K_6\right] \quad (28)$$

Where

$$a_1 = u_0 - a_2 - \frac{A}{M}$$

$$a_2 = \sinh^{-1}\left[\left(u_0 - \frac{A}{M}\right)e^{\sqrt{M}d} + A + (B-A)d\left[\frac{1}{M} + \frac{1}{K}\right] - a_5e^{\sqrt{K}d} - a_6e^{-\sqrt{K}d}\right]$$

$$a_3 = -\left[a_4 + \frac{1}{M}\left[\frac{(a_2^2 - a_1^2)}{4} - \frac{2a_2}{M\sqrt{M}}(B-A) - \frac{2a_1}{M\sqrt{M}}(B-A) - K_2\right]\right]$$

$$a_4 = e^{\sqrt{M}d}\left[a_2e^{\sqrt{K}d} + a_2e^{-\sqrt{K}d} + \frac{1}{K}(K_6 - K_8) - \frac{K_7}{M}\right]$$

$$a_5 = \frac{e^{-\sqrt{K}d}}{\sqrt{K}}[K_{15}]$$

$$a_6 = e^{\sqrt{K}d}[K_{16}]$$

$$a_7 = \left[ \frac{K_9}{K} - a_8 e^{-\sqrt{K}} \right] e^{-\sqrt{K}}$$

$$a_8 = \frac{e^{\sqrt{K}d}}{\sqrt{K}} [K_{10} - K_{11}]$$

$$K_1 = K_{14} - K_{12}$$

$$K_2 = \frac{-a_1^2 + a_2^2}{4} + \frac{2(B-A)}{M\sqrt{M}}(a_2 - a_1)$$

$$K_5 = K_9 - K_6$$

$$K_6 = K_7 + K_8$$

$$K_7 = \frac{-a_1^2 e^{2\sqrt{M}d}}{4} + \frac{a_1^2 e^{-2\sqrt{M}d}}{4} - \frac{(B-A)^2}{2M^2} d^2 + a_1 a_2 M d^2 + \frac{2a_2 e^{-\sqrt{M}d}}{M\sqrt{M}}(B-A) - \frac{2a_1 e^{\sqrt{M}d}}{M\sqrt{M}}(B-A) - K_1 d - K_2$$

$$K_8 = -\frac{a_5^2 K_3 e^{2\sqrt{K}d}}{4K} - \frac{a_6^2 K_3 e^{-2\sqrt{K}d}}{4K} - \frac{K_4}{K^2} d - a_5 a_6 K_3 d^2 - \frac{A + (B-A)d (2(a_5 + a_6)) e^{-\sqrt{K}d}}{Da} - \frac{(2(a_5 + a_6) e^{\sqrt{K}d})(B-A)}{(\sqrt{K})K} - K_5 d$$

$$K_9 = \frac{1}{K} \left[ -\frac{a_5^2 K_3 e^{2\sqrt{K}}}{4K} - \frac{a_6^2 K_3 e^{-2\sqrt{K}}}{4K} - \frac{K_4}{K^2} - a_5 a_6 K_3 - \left( \frac{A + (B-A)d (2(a_5 + a_6)) e^{-\sqrt{K}}}{Da} - \frac{(2(a_5 + a_6) e^{\sqrt{K}})(B-A)}{(\sqrt{K})K} - K_5 + (K_6) \right) \right]$$

$$K_{10} = \frac{1}{K} \left[ -\frac{2\sqrt{K} a_5^2 K_3 e^{2\sqrt{K}d}}{4K} - \frac{2\sqrt{K} a_6^2 K_3 e^{-2\sqrt{K}d}}{4K} - \frac{K_4}{K^2} - 2a_5 a_6 K_3 d - \frac{A + (B-A)d (2(a_5 + a_6)) e^{-\sqrt{K}d}}{Da} - \frac{(B-A) 2(a_5 + a_6) e^{\sqrt{K}d}}{K\sqrt{K}} - K_5 \right]$$

$$K_{11} = \frac{1}{M} \left[ -\frac{2\sqrt{M} a_1^2 e^{2\sqrt{M}d}}{4} - \frac{2\sqrt{M} a_2^2 e^{-2\sqrt{M}d}}{4} - \frac{(B-A)^2}{M^2} d + 2a_1 a_2 M d - \frac{(B-A) 2a_2 e^{-\sqrt{M}d}}{M^2} - \frac{(B-A) 2a_1 e^{\sqrt{M}d}}{M^2} - K_1 \right]$$

$$K_{12} = \sqrt{M} a_3 e^{\sqrt{M}d} - \sqrt{M} a_4 e^{-\sqrt{M}d} - \sqrt{K} a_7 e^{\sqrt{K}d}$$

$$K_{13} = K_{11} + K_{12}$$

$$K_{14} = [MK_{11}]$$

$$K_{15} = \left[ \sqrt{K} a_6 e^{-\sqrt{K}d} + \sqrt{M} (a_1 e^{\sqrt{M}d} - a_2 e^{-\sqrt{M}d}) + (A + (B-A)) \left[ \frac{1}{M} - \frac{1}{K} \right] \right]$$

$$K_{16} = \left[ u_0 - a_5 e^{\sqrt{K}} - \frac{A + (B-A)}{K} \right]$$

## V. RESULTS AND ANALYSIS

To discuss the physical significance of important parameters of flow, 'MATLAB' has been used.

Assuming that the infinite vertical channel moving in a opposite direction, the partly loaded porous medium has been analyzed in order to null over the physical configuration of the laminar Mhd fully developed flow. The momentum and energy equation solutions are obtained by adopting perturbation technique, and the findings are

presented graphically for principle parameters of the flow region. The zero order and first order velocity and temperature distribution are presented for different cases. Due to the linearity, the zero-order temperature profile is not considered for discussion.

The clear fluid at  $y=0$ , interface fluid at region near  $y=d$  and porous region at  $y=1$  have been illustrated.

Case (i) The plate  $y=0$  is heated, when  $A=1$  and the



plate  $y=1$  is cooled, when  $B=0$ .

Case(ii) The plate  $y=0$  is cooled,when  $A=1$  and  $y=1$  is heated, when  $B=1$ .

Fig 1.1 to Fig 1.3 depicts the velocity profile ( $U_{of}$  or  $U_{op}$ ),for case (i) with different parameters Darcy number ( $Da$ ),Magnetic field parameter ( $M$ ) and width ( $d$ ).It is observed that the velocity profile is decreased with the upgrading values of  $Da$  and  $M$ . It is noted that the graph degrades with increasing  $d$ .

Fig 1.4 to fig. 1.6 illustrates the velocity distribution profile ( $U_{of}$  or  $U_{op}$ ) for case(ii).The profile seems to decrease for incrementing values of  $Da$  and  $M$  but the profile is elevated for ascending values of  $d$ .

From fig 2.1 to fig.2.6 it is observed that the velocity ( $U_{lf}$  or  $U_{lp}$ ) flow proliferates with increasing  $M$  and  $d$  for either  $A$  or  $B$  exists. But the Darcy number  $Da$  decelerates the velocity flow for case (i) and it accelerates the flow for case (ii)

From fig 3.1 to fig. 3.6 it is observed that the effect of the parameters  $M$  and  $Da$ . It is noted that the temperature distribution of the flow enhances for either cases. But for increasing  $d$  values, it is noted that the temperature profile decelerates for case (i) and accelerates for case (ii).

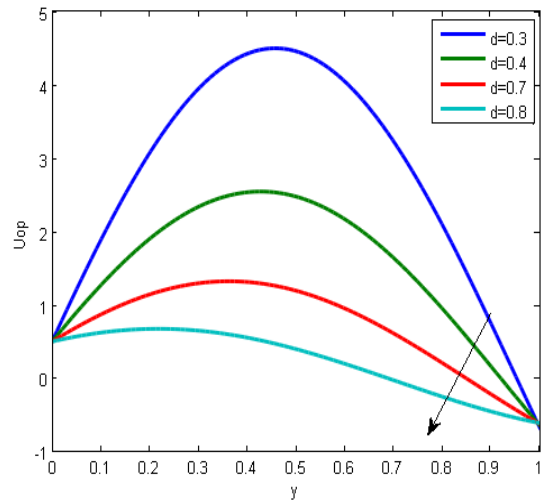


Fig 1.3: Velocity stretch for  $d$ : [ $Da=0.1$ ;  $A=1$ ;  $B=0$ ;  $u_0=0.5$ ;  $M=2$ ]

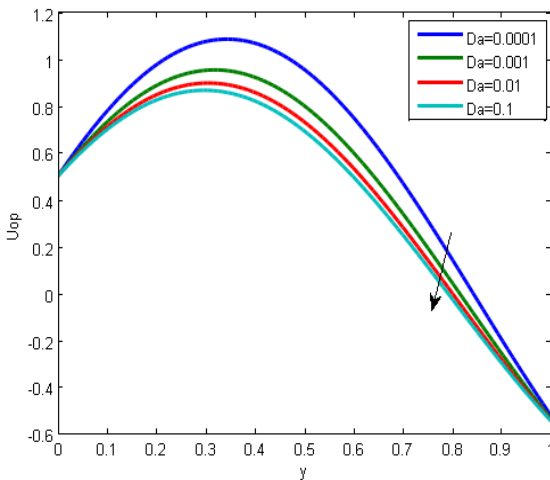


Fig 1.1: Velocity profile for  $Da$ , [ $A=1$ ;  $B=0$ ;  $u_0=0.5$ ;  $d=0.3$ ;  $M=2$ ]

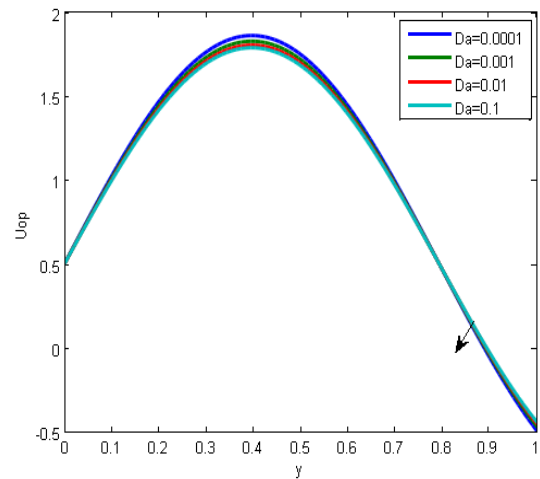


Fig 1.4: Velocity profile for  $Da$ : [ $A=0$ ;  $B=1$ ;  $u_0=0.5$ ;  $d=0.3$ ;  $M=2$ ]

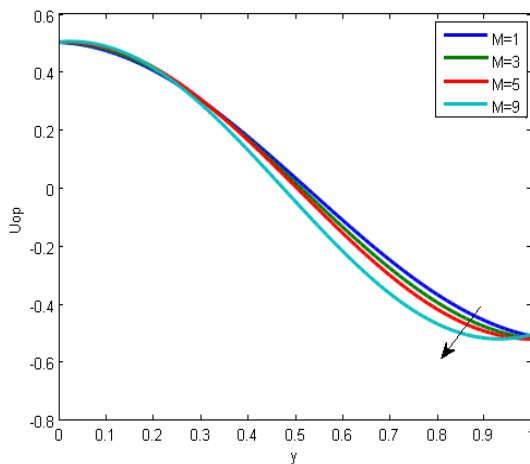


Fig 1.2: Velocity profile for  $M$ , [ $A=1$ ;  $B=0$ ;  $u_0=0$ ;  $d=0.3$ ;  $Da=0.1$ ]

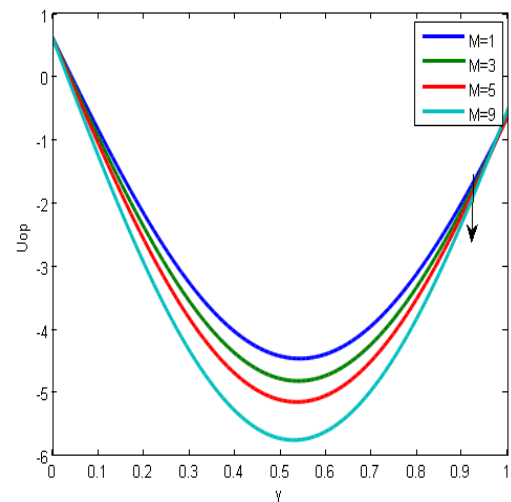


Fig 1.5: Graph of velocity  $U_{op}$  for  $M$ : [ $Da=0.1$ ,  $A=0$ ;  $B=1$ ;  $u_0=0.5$ ;  $d=0.3$ ]



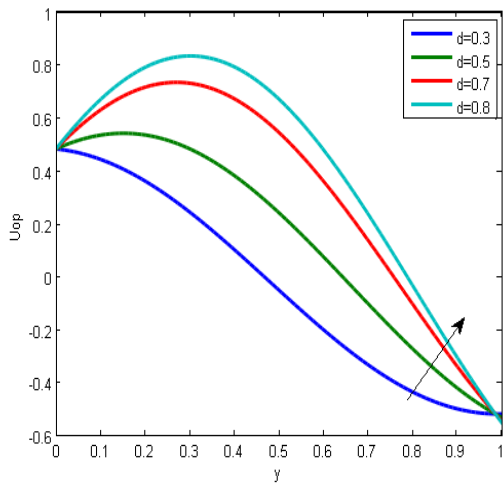


Fig 1.6: Graph of velocity  $U_{op}$  for  $d$ ; [ $Da=0.1$ ,  $A=0; B=1; u_0=0.5; M=2$ ]

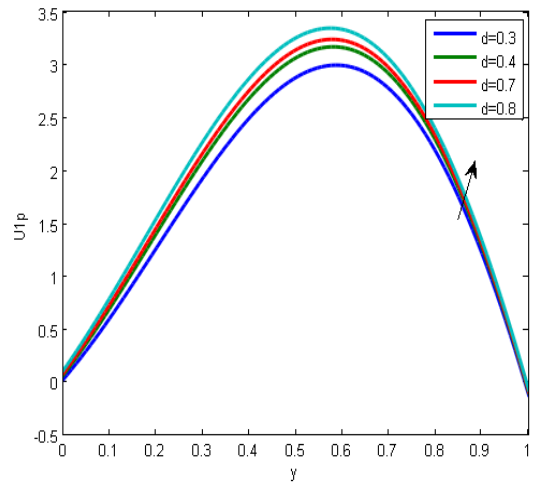


Fig 2.3: Graph of velocity  $U_{1p}$  for  $d$ ; [ $M=2$ ,  $A=1; B=0; u_0=0.5; Da=0.1$ ]

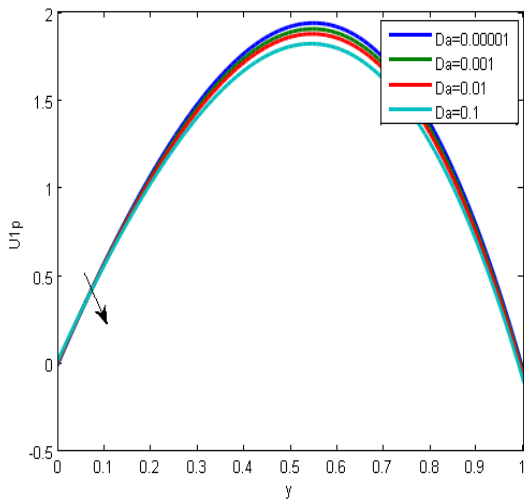


Fig 2.1: Graph of velocity  $U_{1p}$  for  $Da$ ; [ $M=2$ ,  $A=1; B=0; u_0=0.5; d=0.3$ ]

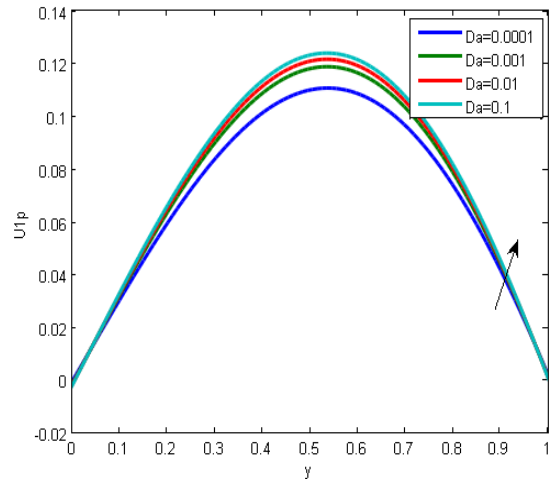


Fig 2.4: Graph of velocity  $U_{1p}$  for  $Da$ ; [ $M=2, A=0; B=1; u_0=0.5; d=0.3$ ]

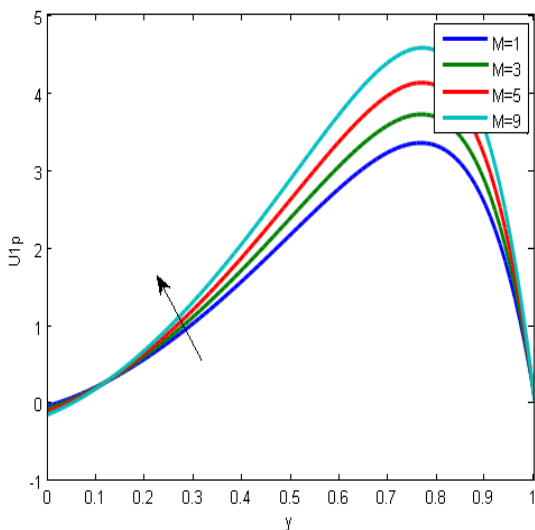


Fig 2.2: Profile of velocity  $U_{1p}$  for  $M$ ; [ $Da=0.1$ ,  $A=1; B=0; u_0=0.5; d=0.3$ ]

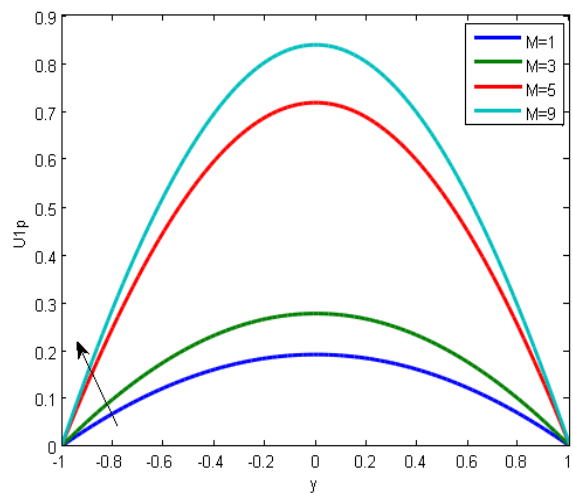


Fig 2.5: Profile of velocity  $U_{1p}$  for  $M$ ; [ $Da=0.1$ ,  $A=0; B=1; u_0=0.5; d=0.3$ ]

The Influence of Inclined Magnetic Field and Heat Transfer on the MHD Convective Flow in a Vertical Channel Filled Partly with Porous Medium

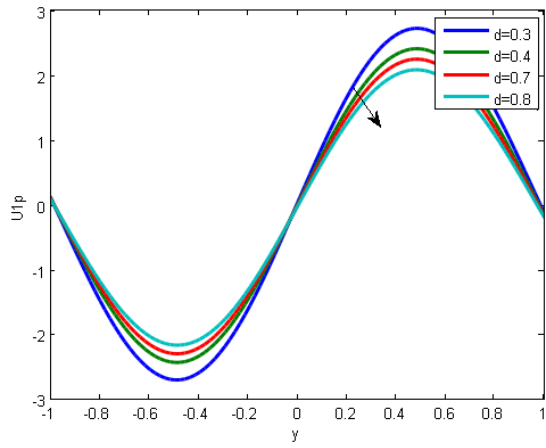


Fig 2.6: stretch of velocity  $U1p$  for  $d$ ; [ $M=2, A=0; B=1; u_0=0.5; Da=0.1$ ]

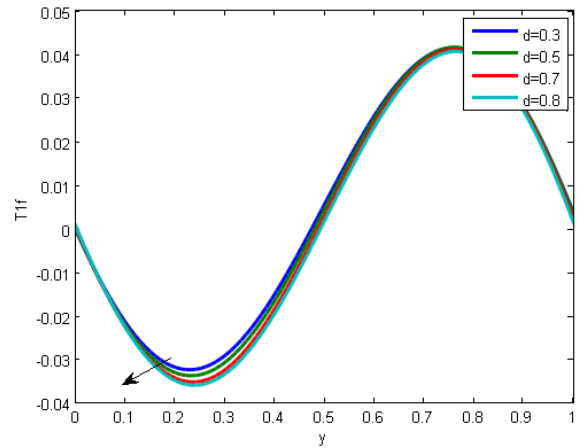


Fig 3.3: Stretch of temperature  $T1f$  for  $d$ ; [ $M=2, A=1; B=0; u_0=0.5; Da=0.1$ ]

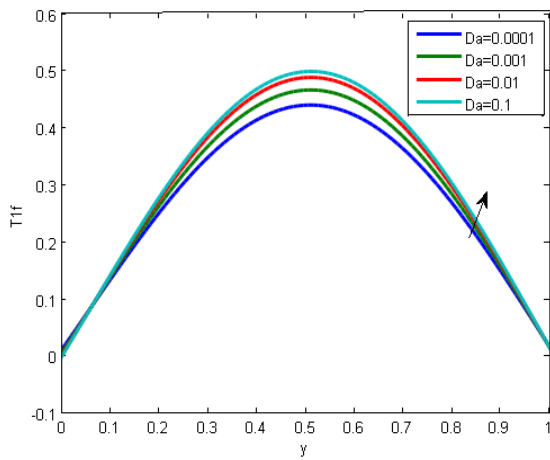


Fig 3.1: Graph of temperature  $T1f$  for  $Da$ ; [ $M=2, A=1; B=0; u_0=0.5; d=0.3$ ]

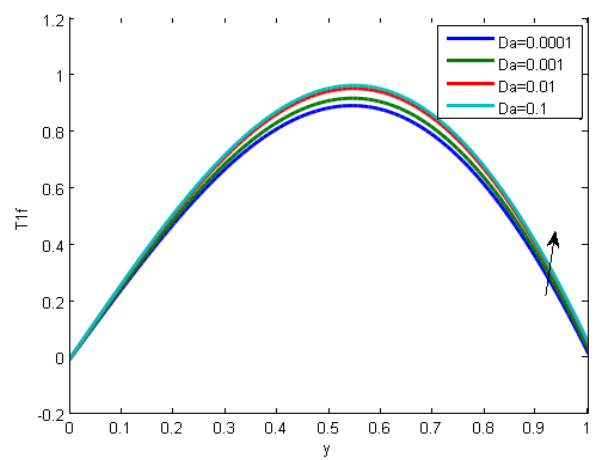


Fig 3.4: Temperature profile  $T1f$  for  $Da$ ; [ $M=2, A=0; B=1; u_0=0.5; d=0.3$ ]

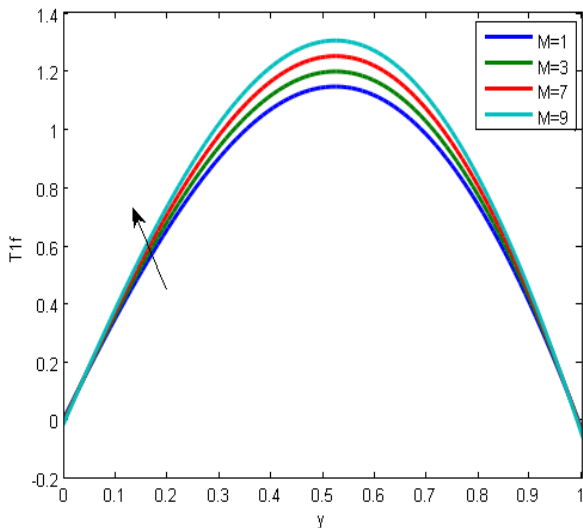


Fig 3.2: Graph of temperature  $T1f$  for  $M$ ; [ $M=2, A=1; B=0; u_0=0.5; Da=0.1$ ]

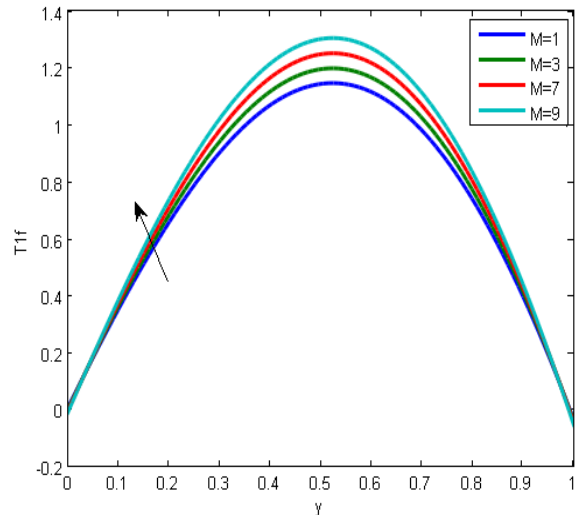


Fig 3.5: Stretch of temperature  $T1f$  for  $M$ ; [ $d=0.3, A=0; B=1; u_0=0.5; Da=0.1$ ]



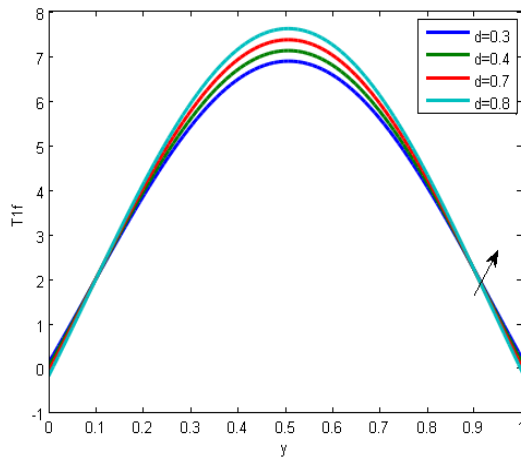


Fig 3.6: Graph of temperature  $T_{1f}$  for  $d$ : [ $M=2$ ,  $A=0$ ;  $B=1$ ;  $u_0=0.5$ ;  $Da=0.1$ ]

## VI. DISCUSSION AND CONCLUSION

The analysis of a laminar fully developed Mhd flow bounded by vertically parallel plates which are moving in opposite direction and maintained at different temperature in presence of inclined magnetic field has been studied. The expression for velocity and temperature are obtained by implementing perturbation technique. It is imperative to conclude that for case (i) the velocity profile ( $U_{of}$  or  $U_{op}$ ) degrades for incrementing values of parameters  $d$ ,  $M$  and  $Da$ . Whereas in case (ii) the velocity profile shows reduction with inclusion of  $M$  and  $Da$  but it accelerates with the increasing  $d$  value. On the other hand, the velocity profile ( $U_{lf}$  or  $U_{lp}$ ) enhances the flow for case (i) with variation of  $M$ ,  $d$  and it seems to decrease for increasing  $Da$  value. In case (ii) the velocity profile progresses for ascending values of  $M$  and  $Da$  but it decreases for increasing  $d$  values. As per the temperature profile is concerned, the profile ( $\theta_{lf}$  or  $\theta_{lp}$ ) for case (i) it seems to decrease by incrementing  $d$  values and increases for increasing magnetic field  $M$  and darcy number  $Da$ . It is noted from case (ii) that the temperature profile communicates the same increasing behavior with the predominant parameters  $Da$ ,  $M$  and  $d$ .

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